
Ordering effects in electric splay Freedericksz transitions

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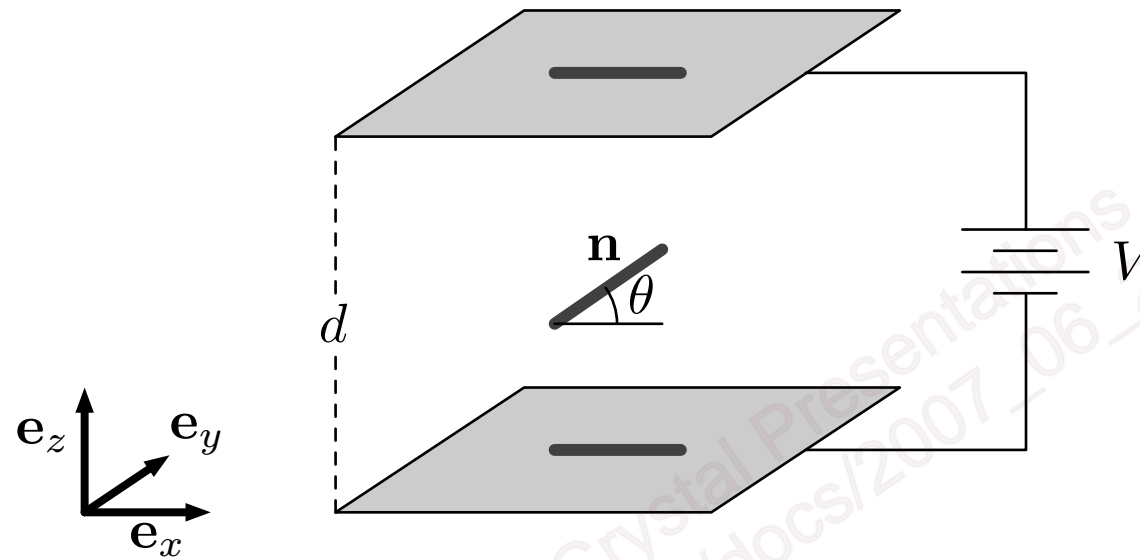
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Motivation

We aim at investigating the effects of a variable degree of orientation on the equilibrium configurations in electric splay Fredericksz geometry

- Small voltages *disorder* the sample
- Dependence of V_{cr} on s_0
- Most interesting regime: high anisotropy ($\epsilon_{\perp} \ll \epsilon_{\parallel}$)

Geometry



Dimensionless variables $u = z/d \in [0, 1]$, $\phi = \varphi/V$, such that

$$\mathbf{E} = \nabla \varphi = \frac{V}{d} \phi'(u) \mathbf{e}_z$$

Uniaxial order tensor $\mathbf{Q}(u) = s(u) (\mathbf{n} \otimes \mathbf{n} - \frac{1}{3} \mathbf{I})$, with director

$$\mathbf{n}(u) = \cos \theta(u) \mathbf{e}_x + \sin \theta(u) \mathbf{e}_z$$

Free energy density

$$\Psi = \Psi_{\text{el}}(\mathbf{Q}, \nabla \mathbf{Q}) + \Psi_E(\mathbf{Q}) + \Psi_{\text{LdG}}(\mathbf{Q})$$

with

$$\Psi_{\text{el}} = \frac{1}{2} K |\nabla \mathbf{Q}|^2 = K \left(s^2 |\nabla \mathbf{n}|^2 + \frac{1}{3} |\nabla s|^2 \right)$$

$$\Psi_E = -\frac{1}{2} \mathbf{D} \cdot \mathbf{E} = -\frac{1}{2} (\bar{\epsilon} \mathbf{E} + \epsilon_a \mathbf{Q} \mathbf{E}) \cdot \mathbf{E}$$

$$\Psi_{\text{LdG}} = A \text{tr} \mathbf{Q}^2 - B \text{tr} \mathbf{Q}^3 + C \text{tr} \mathbf{Q}^4 = \frac{2}{3} A s^2 - \frac{2}{9} B s^3 + \frac{2}{9} C s^4$$

where

$$\bar{\epsilon} = \frac{1}{3} (\epsilon_{\parallel} + 2\epsilon_{\perp}), \quad \epsilon_a = \epsilon_{\parallel} - \epsilon_{\perp} > 0, \quad \gamma_a = \frac{\epsilon_a}{3\bar{\epsilon}} \in [0, 1)$$

Characteristic lengths

$$\xi_e^2 = \frac{2Kd^2}{\epsilon_a V^2}, \quad \xi_n^2 = \frac{9K}{4C} = 6\eta^2 \xi_e^2 \quad (\eta \lesssim 1)$$

Equilibrium equations

Euler-Lagrange + Maxwell equations

$$s\theta'' + 2s'\theta' + \frac{d^2}{\xi_e^2} \phi'^2 \sin \theta \cos \theta = 0$$

$$\frac{2}{3}s'' - 2s\theta'^2 - \frac{d^2}{\xi_n^2} \frac{d\sigma}{ds} + \frac{d^2}{\xi_e^2} \phi'^2 \left(\sin^2 \theta - \frac{1}{3} \right) = 0$$

$$\frac{d}{du} \left(\phi'(1 - \gamma_a s) + 3\gamma_a s \phi' \sin^2 \theta \right) = 0$$

Boundary conditions

$$\theta(0) = \theta(1) = 0, \quad s'(0) = s'(1) = 0, \quad \phi(0) = 0, \quad \phi(1) = 1$$

Results: Freedericksz transition

- Critical voltage $V_{\text{cr}} = \sqrt{\frac{2\pi^2 K s_0}{\epsilon_a}}$
(renormalization of elastic constant)

- Shape of minimizers. Let $V = V_{\text{cr}} (1 + \epsilon^2)$

$$\theta(u) = \epsilon A_1 \sin \pi u + o(\epsilon^2)$$

$$s(u) = s_0 - \epsilon^2 \eta^2 (B_1 + B_2 \cos 2\pi u) + o(\epsilon^2)$$

$$\phi(u) = u + C_1 \epsilon^2 \sin 2\pi u + o(\epsilon^2)$$

with positive integration constants and, in particular,

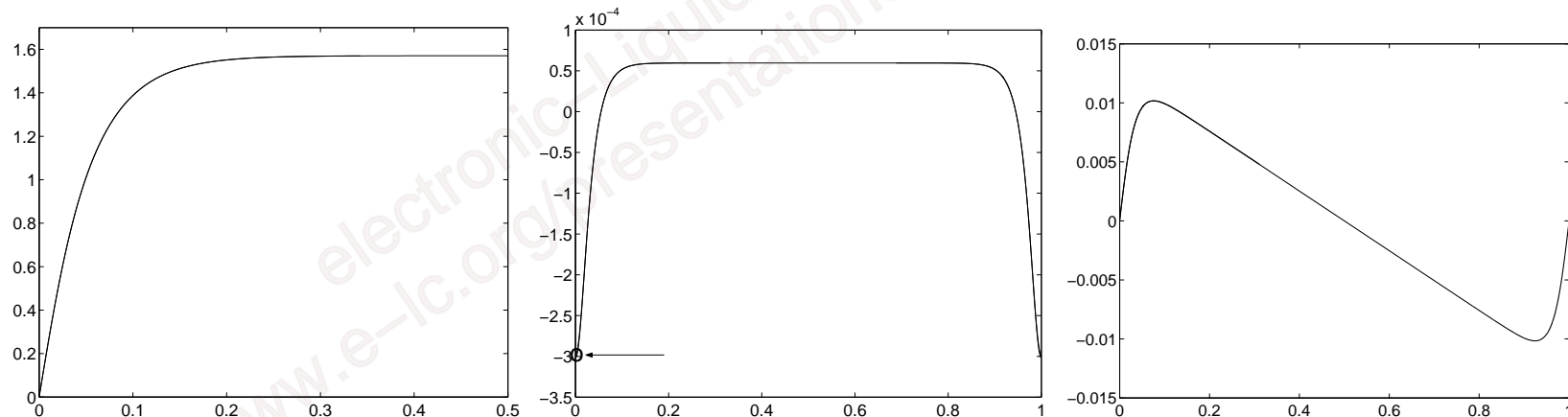
$$A_1^2 = \frac{4(1 - \gamma_a s_0)}{1 + 2\gamma_a s_0} + O(\eta^2)$$

Results: Strong applied voltages

In the limit $V \gg V_{cr}$ the tilt angle develops a boundary layer to maximize the size of the region where the director is parallel to the external field.

The strong director gradient induce a surface melting

Nematic effects decrease the electric field in the bulk



High anisotropy

Let $\gamma_a = 1 - \epsilon^2$ ($\epsilon_{\perp}/\epsilon_{\parallel} = O(\epsilon^2)$)

Singular in Frank limit (Self, Please, Sluckin, 2002)

- Close to the Fredericksz transition the leading orders of the relevant fields are modified: both θ and ϕ become $O(\epsilon^2)$
- For strong applied voltages a double boundary-layer structure arises, and the electric potential develops a finite jump in an infinitesimal layer

Regularized by a finite degree of orientation

- Close to the Fredericksz transition the leading orders are unchanged: $\theta = O(\epsilon)$ while $\phi = O(\epsilon^2)$
- Even for strong fields the electric field remains bounded

Conclusions

- Nontrivial disordered regions/regimes in the presence of an external field
- Regularizing role of a finite degree of orientation in the high anisotropy limit
- From the mathematical point of view, regularization is related to the gain of coercivity in the free energy functional

Comments? Reprints?

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